

Rotational Structures of Neutron Deficient Isotopes $^{125,127,129}\text{Ce}^*$

MA Hai-Liang¹⁾ DONG Bao-Guo

(China Institute of Atomic Energy, Beijing 102413, China)

Abstract The configuration-dependent cranked Nilsson-Strutinsky approach has been used to investigate rotational structures of neutron deficient isotopes $^{125,127,129}\text{Ce}$. Signature splitting and deformation of yrast bands have been discussed. Shape coexistence may exist in $^{127,129}\text{Ce}$. The signature splitting of yrast bands for Ce isotopes is strongly dependent on occupied orbitals because of slight triaxial deformation.

Key words signature splitting, shape coexistence, triaxial deformation

1 Introduction

$^{125,127,129}\text{Ce}$ isotopes have been populated experimentally to high spin states in past decades^[1–4]. These isotopes are well deformed with quadrupole deformation of $\beta_2 = 0.25 - 0.30$ ^[5]. In present paper, level schemes, shape coexistence and signature splitting have been studied in the framework of configuration-dependent cranked Nilsson-Strutinsky (CNS) approach^[6, 7]. The Nilsson parameters are quoted from Ref. [6] and pairing correlations neglected, so it is qualitatively to compare the calculated bands with observed bands in the low spin region $I < 15\hbar$. In the high spin region, where pairing correlation will play less important role or is negligible, the discrepancy between experimental results and cranked shell model calculations will become small^[8].

In the following a short notation will be used to label the configurations (relative to $^{100}\text{Sn}_{50}$ core), $[p, n] \equiv \pi(h_{11/2})^p \otimes \nu(h_{11/2})^n$. The simplification for signature is used: $(\alpha = +) \equiv (\alpha = +1/2)$ and $(\alpha = -) \equiv (\alpha = -1/2)$. Then we have four combinations of parity and signature indicated by (π, α) .

The experimental and theoretical $E - E_{\text{RLD}}$ energies as a function of spin are shown in Fig. 1, where

$E_{\text{RLD}} \equiv (\hbar^2/2J_{\text{rig}})I(I+1) \equiv 32.32A^{-5/3}I(I+1)$ MeV is the energy of a rigid rotor reference with A being the mass number. Here we only show the configurations near yrast line and left panels for $\pi = +$ and right panels $\pi = -$. In the figure, the experimental and theoretical $E - E_{\text{RLD}}$ values have been renormalized at the highest spin state for both parities.

2 Signature splitting

Based on the CNS calculations, the [2,7] configuration is assigned to the negative parity yrast bands. It is the same as previous works^[1–4]. For the positive parity yrast bands, its configuration is assigned as [2,6] for ^{125}Ce and [2,8] for ^{129}Ce . For positive parity band of ^{127}Ce , it was observed only up to $25/2\hbar$, so we do not show experimental result in Fig. 1.

Signature splitting is the energy shift between energetically favored and unfavored bands. The energy staggerings of the calculated [2,6] configuration and the observed band 3 in ^{125}Ce are shown in Fig. 2. Here we use the same definition as in Ref. [9],

$$\Delta E(I) = \frac{E(I) - E(I-1) - E(I+1) + E(I) + E(I-1) - E(I-2)}{2}.$$

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1) E-mail: mhl624@iris.ciae.ac.cn

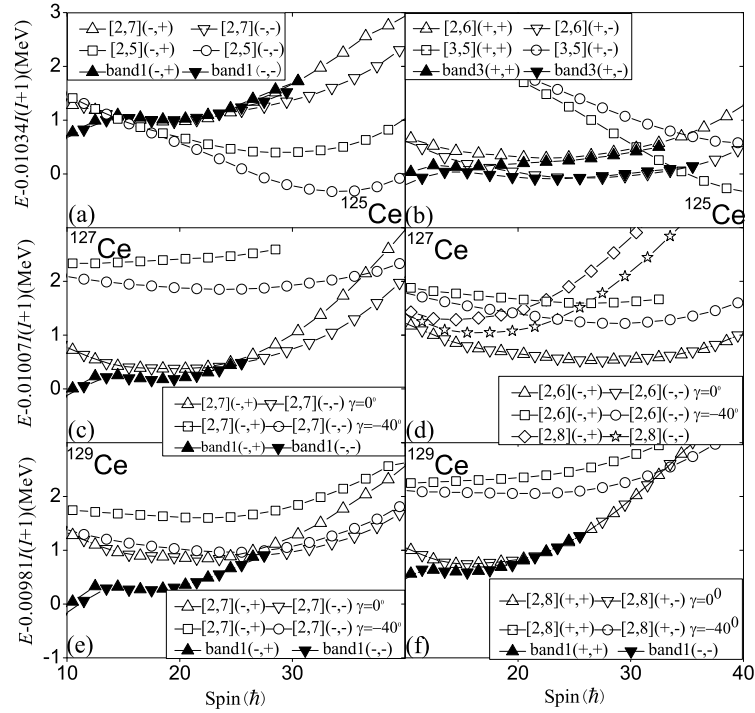


Fig. 1. Energies as a function of spin for the calculated and observed bands of interest for $^{125,127,129}\text{Ce}$. The open symbols indicate the calculated data and the solid ones the observed data.

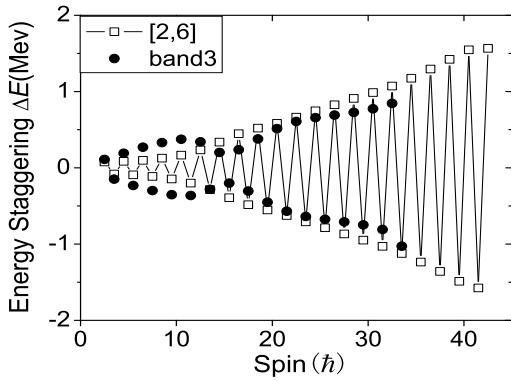


Fig. 2. Energy staggerings $\Delta E(I)$ plotted versus spin for the calculated [2,6] configuration and observed band 3 in ^{125}Ce .

The calculated shape trajectories as a function of spin corresponding to the calculated yrast bands are shown in Fig. 3. It can be seen from Fig. 3 that the deformation of configuration [2,6] is near axial. So even if no triaxial deformation, large signature splitting still can exist.

For $^{127,129}\text{Ce}$, very small signature splitting will be expected for the yrast bands since it can be seen in Fig. 1 that the calculated and observed signature partner bands are nearly degenerated. It is confirmed

that the CNS approach reproduces signature splittings well in the high spin region.

3 Nuclear shapes

In Ref. [10], the γ driving force of odd-odd Pr isotopes has been studied. Since $^{125,127,129}\text{Ce}$ isotopes have the same neutron number as $^{126,128,130}\text{Pr}$, the γ driving force is similar. In our CNS calculations, all yrast bands show very small values of triaxial deformation, $\gamma \sim 0^\circ$, except the negative parity configuration [2,7] of ^{125}Ce . The calculated potential energy surfaces (PES) for this negative parity [2,7] configuration in ^{125}Ce show that this configuration is very γ -soft.

There are shape coexistences within the same configuration in $^{127,129}\text{Ce}$ which means there are two minima in the PES. One is located at around $\gamma \sim 0^\circ$ and the other is located at around $\gamma \sim -40^\circ$. Both bands built on the two minima are shown in Fig. 1. The $\gamma \sim -40^\circ$ bands are higher in energy than the $\gamma \sim 0^\circ$ bands, the favored bands, in the interesting spin region. According to this band structures, observed yrast bands only have slight triaxial deformation in $^{127,129}\text{Ce}$.

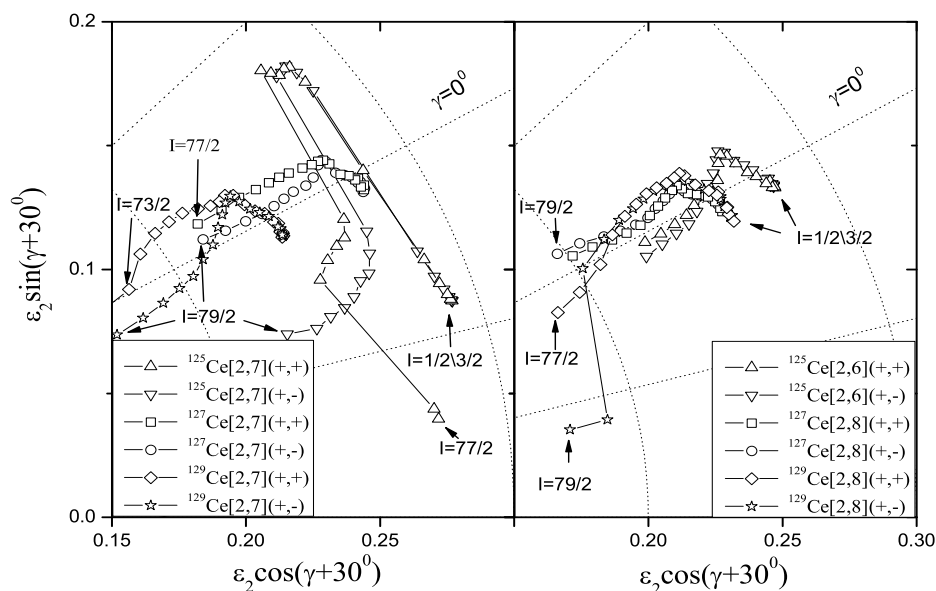


Fig. 3. Calculated shape trajectories of the yrast configurations as a function of spin in the (ϵ_2, γ) -plane for $^{125,127,129}\text{Ce}$.

4 Conclusions

The band structures, nuclear shape and signature splitting of neutron deficient odd-even cerium isotopes $^{125,127,129}\text{Ce}$ have been investigated in the framework of CNS approach. The triaxial deforma-

tion of yrast bands in $^{125,127,129}\text{Ce}$ is small, $\gamma \sim 0^\circ$. There is shape coexistence in $^{127,129}\text{Ce}$ within the same configuration. Even no triaxial deformation, large signature splitting can exist. Good agreement is found between theoretical and experimental results when spin is over $15\hbar$ for these Ce isotopes.

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缺中子核素 $^{125,127,129}\text{Ce}$ 的转动结构*

马海亮¹⁾ 董保国

(中国原子能科学研究院 北京 102413)

摘要 用组态相关推转 Nilsson-Strutinsky 方法研究了缺中子核素 $^{125,127,129}\text{Ce}$ 的转动结构性质, 讨论了转晕带的旋称劈裂和核的形状. 转晕带相应的核形状均是近轴对称的. $^{127,129}\text{Ce}$ 在同一个组态中可能有形状共存. 理论指出即使对于无三轴形变核也有可能出现较大的旋称劈裂.

关键词 旋称劈裂 形状共存 三轴形变

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1) E-mail: mhl624@iris.ciae.ac.cn